





Thermomechanical Measurements for Energy Systems (MENR)

Measurements for Mechanical Systems and Production (MMER)

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Engines (M) that produce mechanical energy are always coupled with an **Operating Machine** (U) through a <u>shaft</u>, that transfers the energy by an alternating (v) or a rotating (ω) movement.

Relative movements of mechanical parts of engines and operating machines always produce spurious *vibration* and audible *noise*.

Vibrations are a "mechanical energy waste", are troublesome for operators and harmful to machines, because of the "mechanical fatigue effects" they cause !

example of an agricultural tractor engine:



Vibrations are <u>harmonic motions of mechanical parts</u> of a machine which, sometimes, can be seen by eye and always, can be heard as noise ...



Displacement:
$$x(t) = Xsen\omega t = Xe^{j\omega t}$$
Velocity: $v(t) = X\omega \cos \omega t = X\omega e^{j\omega t}$ Acceleration: $a(t) = -X\omega^2 sen\omega t = -X\omega^2 e^{j\omega t}$

Therefore, **vibrations** are characterized by **waveforms** which have their specific **amplitude** and **frequency** !

Every *periodic waveform* can be described by a *Fourier series* :

$$x(t) = x(t + 2\pi) = \frac{1}{2}a_0 + \sum_{k=1}^{\infty} \left(a_k \cos kt + b_k \operatorname{senkt}\right)$$

therefore, to study vibrations we can refer to its component waves

As for a mechanical system in vibration, the momentum $Q = m \cdot v = \cos t$ remains constant, these three simple relations indicate that it is preferable to **measure displacements** at *low frequencies* and to **measure acceleration** at *high-frequency* !!



The main problem for *vibration measurements* is finding a *fixed point* that can be considered as *reference point* for all the points or elements of the "vibrating machine" !





O – oculare con reticolo

Sometimes, this can be done placing a *telescope instrument* equipped with a *graduated aim* quite far form the vibrating machine !

However, most of the times in industrial plants, machines are placed indoor in a shed !

The *fixed point*, therefore, can *NOT BE FAR* from the vibrating machine !



A <u>seismic mass</u> m is placed inside the transducer, connected to the case by means of a <u>spring</u> k and a <u>viscous damper</u> c ...



Seismic transducers utilize an "internal fixed reference point" and are *rigidly attached* to the engine !

How are these sensors actually working ... ?



We would like to measure the **periodic displacement** $u(t) = U_0 sen\omega t$ of the machine M, indicated in the figure on the <u>imaginary</u> absolute (fixed) scale !

However, the transducer output is the **displacement** $\delta(t)$ of the *inner seismic mass* «*m*» with respect to the sensor external case, indicated in the figure on the <u>real</u> relative (moving) scale ...

$$u(t) \longrightarrow \text{TRSD} \longrightarrow \delta(t)$$

We can express the motion of the seismic mass m by means of the theorem of relative motion

 $x(t) = u(t) + \delta(t)$

- **x(t)** absolute displacement (measured in an inertial reference frame)
- *u(t) dragging displacement* (of the relative reference frame)
- $\delta(t)$ relative displacement (measured in the relative reference frame)

Vibrations are definitely a dynamic phenomenon so, we will study first the dynamic response of the seismic transducer: that means finding the gain δ_0/U_0 with respect to the frequency ω of the displacement u(t) ...

Seismic transducers are clearly *second order instruments*, so the equation for the dynamic response is :

 $m\frac{d^2x}{dt^2} + c\frac{d\delta}{dt} + k\delta = 0$

because <u>inertial forces</u> are always referred to the *inertial reference frame* (x), while <u>elastic and damping forces</u> here are referred to the *relative moving frame* (δ) made of the sensor case !

 $m\frac{d^2}{dt^2}(\delta+u) + c\frac{d\delta}{dt} + k\delta = 0$

$$m\frac{d^{2}\delta}{dt^{2}} + c\frac{d\delta}{dt} + k\delta = -m\frac{d^{2}u}{dt^{2}} = -m(-\omega^{2}U_{o}sen\omega t) = m\omega^{2}u(t)$$

the "external force term" in the equation of this transducer is proportional to the relative displacement u(t) ... that is $F(t) = m\omega^2 U_0 \operatorname{sen}\omega t$

Frequency response is much similar to any other *second order instrument* :





Seismic-displacement-pickup frequency response.



which is the <u>frequency response</u> of the relative displacement $\delta(t)$ of the seismic mass m with respect to the transducer case !

it can be seen on the curves that for every **damping factor** $\boldsymbol{\zeta}$, the usable frequency range is for $\boldsymbol{\omega} >> \boldsymbol{\omega}_n$ where the <u>gain</u> is $\boldsymbol{\delta}_0 / \boldsymbol{U}_0 \approx \boldsymbol{1}$ and the <u>phase delay</u> is $\boldsymbol{\varphi} \approx -\boldsymbol{180}^\circ \dots$ in *phase opposition*

these conditions together imply that when the machine oscillation u(t) is at its maximum point U_0 the relative displacement $\delta(t)$ of the seismic mass m is at its minimum point $-\delta_0$...

This means the *seismic mass m* actually «stopped moving» with respect to the absolute reference frame: in these conditions, the *seismic mass m* realizes the *fixed reference point* for vibration measurements !!

Operative considerations:

- To make the condition $\omega \gg \omega_n$ really usable, we need to design a transducer with a very low $\omega_n = \sqrt{\frac{k}{m}}$ which implies $\underline{k} \rightarrow very small}$ and $\underline{m} \rightarrow very big}$!
- It is quite hard to simultaneously realize these two conditions (a very soft spring that bears a big mass !!)
- Moreover, a big seismic mass m firmly applied on the machine mass M causes a «change of the vibration modes» and, therefore, induces an insertion error !

Few transducers realize the conditions of above : the **vibrometer** and the **seismograph** ...

In practice, it would be much better to make measurements with a transducer made with a «small mass m» and a «high stiffness k»

This implies a **high natural frequency** $\omega_n = \sqrt{\frac{k}{m}}$ and making measurement at frequencies $\omega \ll \omega_n$

Because the gain δ_o/U_o is quite low for $\omega \ll \omega_n$ we can not directly measure the *displacement* u(t) in these conditions and we have to *change the measurand* at the transducer input !

This implies to change transducer type ...

In fact, being $\ddot{u}(t) = -\omega^2 \cdot U_0 sen\omega t$ for the

which substituted in the *frequency response* of the vibrometer results :

$$\delta_{o} = U_{o} \frac{\omega^{2}}{\omega_{n}^{2}} \cdot \frac{1}{\sqrt{\left(1 - \frac{\omega^{2}}{\omega_{n}^{2}}\right)^{2} + 4\xi^{2} \left(\frac{\omega}{\omega_{n}}\right)^{2}}} = \frac{\ddot{U}_{o}}{\omega^{2}} \cdot \frac{\omega^{2}}{\omega_{n}^{2}} \cdot \frac{\omega^{2}}{\sqrt{\left(1 - \frac{\omega^{2}}{\omega_{n}^{2}}\right)^{2} + 4\xi^{2} \left(\frac{\omega}{\omega_{n}}\right)^{2}}}$$

We need to employ a sensor made with the same constitutive elements of the vibrometer (*m*; *c*; *k*), but measuring the <u>drag acceleration $\ddot{u}(t)$ </u>

Such a transducer is the *accelerometer* which has a *frequency response*

$$\ddot{u}(t) \longrightarrow \text{ACCEL} \longrightarrow \delta(t)$$

$$G = \frac{\delta_o}{\ddot{U}_o} = \frac{1}{\omega_n^2} \cdot \frac{1}{\sqrt{\left(1 - \frac{\omega^2}{\omega_n^2}\right)^2 + 4\xi^2 \left(\frac{\omega}{\omega_n}\right)^2}}$$

$$\ddot{U}_o = \omega^2 U_o \qquad U_o = \frac{U_o}{\omega^2}$$

amplitudes it results :
$$\ddot{U}_o = \omega^2 U$$

$$U_o = \omega^2 U_o$$
 $U_o = \frac{\omega}{\omega}$

e amplitudes it results :
$$\ddot{U}_o = a$$

$$G = \frac{\delta_o}{\ddot{U}_o} = \frac{1}{\omega_n^2} \cdot \frac{1}{\sqrt{\left(1 - \frac{\omega^2}{\omega_n^2}\right)^2 + 4\xi^2 \left(\frac{\omega}{\omega_n}\right)^2}}$$



Piezoelectric accelerometer frequency response

Piezoelectric accelerometers are still the best acceleration transducers available on the market, followed by *capacitive* and *semiconductor accelerometers* ...

The usable frequency range is now at the left of the *resonance frequency* ω_n !

To extend the frequency band we have to increase ω_n however, this implies a decrease

of the gain $\frac{\delta_o}{\ddot{U}_o}$ or the sensor sensitivity !

If we wish to get the displacements u(t) with such a sensor, we have to proceed with a double integration of the accelerometer signal ...

$$\ddot{u}(t) \longrightarrow \textbf{ACCL} \longrightarrow \delta(t) \longrightarrow \iint \delta(t) dt \longrightarrow u(t)$$



(c)





(d)



Piezoelectric accelerometer designs.

- (a) Peripheral-mounted compression design.
- (b) Center-mounted compression design.
- (c) Inverted center-mounted compression design.
- (d) Shear design.







S: Spring M: Mass P: Piezoelectric element B: Base C: Cable (Courtesy Bruel and Kjaer Instruments, Marlboro, Mass.)



 The piezoelectric effect causes crystal materials like quartz to generate an electric charge when the crystal material is compressed, twisted, or pulled. The reverse also is true, as the crystal material compresses or expands when an electric voltage is applied. Although the most common *piezoelectric crystal* is <u>quartz</u>, to realize *piezoelectric transducers*, titanates of barium (BaTiO₃) are generally used, or the so-called <u>*PZT*</u> (from Lead - Zirconium - Titanium, solid solutions of PbZrO₃ and PbTiO₃)







Piezoelectric quartz crystallizes in the hexagonal system. A slice is cut and extracted from the crystal such as in the figure

When applying a *force F* parallel to the <u>y axis</u>, a *charge ±Q* appears on the <u>x faces</u>

The charge stays as long as force is applied on the y faces: $Q = k_{ij} \cdot F$ k_{ij} are the *piezoelectric constants*

Because of the form they are cut, *piezoelectric* crystals can be considered as **capacitors** with a capacity : $C = \varepsilon_0 \varepsilon_r S / d$ The charge then is : $Q = C \cdot V$ and the **graduation curve** is :

$$V = \frac{Q}{C} = \frac{k_{ij} \cdot F}{\varepsilon_0 \varepsilon_r S / d} = \frac{k_{ij} \cdot m \ddot{y}}{\varepsilon_0 \varepsilon_r S / d} = \frac{dk_{ij}m}{\varepsilon_0 \varepsilon_r S} \cdot \ddot{y}$$



m is the *seismic mass* that during vibration presses on the piezoelectric crystal with the *inertial forces* $F = m\ddot{y}$

Note that *m* is mechanically *pre-charged* to the crystal so to output a signal also for negative inertial forces ... *c* and *k* are the crystal internal *damping* and *stiffness* ...



Piezoelectric transducers, unfortunately, have a very high output impedance ($10^9 \Omega$) which makes them very poor signal generators. To actually read the signal proportional to the acceleration \ddot{y} we have to employ special amplifiers, with very high input impedance ($10^{12} \Omega$): the **charge amplifiers** !



Many elements play a role in this delicate *transducer–amplifier coupling* :

- C_q = piezoelectric quartz capacity C_c = leads coupling capacity C_a = amplifier input capacity
 - R_{α} = piezoelectric quartz output resistance $R_c = leads resistance$ R_a = amplifier input resistance

We have all passive elements in parallel :

$$C = C_q + C_c + C_a \qquad \frac{1}{R} = \frac{1}{R_q} + \frac{1}{R_c} + \frac{1}{R_a} \qquad R = \frac{R_q R_c R_a}{R_c R_a + R_q R_a + R_q R_c}$$



It discharges exponentially, like a 1° order electrical system : $V(t) = V_{\alpha}e$

Therefore, charge amplifiers are made with <u>FET input stage inverting AO</u> and a feedback capacitor :







Note that the graduation curve $V = \frac{dk_{ij}m}{\varepsilon_0\varepsilon_rS}$. \ddot{y} was derived

without taking into account that quartz is still an *elastic element* (with a <u>very high</u> Young's modulus E) and that the charge Q becomes available on the faces of the crystal due to an internal electrical asymmetry, caused by the very small deformation ε of the crystal. The deformation ε in turn is caused by the *inertia forces* of the seismic mass m which acts during vibration

$$Y = Q$$

 F
 F
 $Al = y$
 Z

$$Q = k_{ij} \cdot F = k_{ij} \cdot S \sigma = k_{ij} \cdot S \cdot E\varepsilon = k_{ij} \cdot SE \cdot \frac{\Delta l}{l} \qquad \qquad Q = \frac{k_{ij}SE}{l} \cdot y_q$$

On the node A in the figure above, we have: $i_q = -i_f$ $i_q = \frac{dQ}{dt} = K \cdot \dot{y}_q$ being $\frac{k_{ij}SE}{l} = \cos t = K$ $i_f = C_f \frac{dV_o}{dt} = C_f \cdot \dot{V}_o$ therefore: $K \cdot \dot{y}_q = -C_f \cdot \dot{V}_o$

The graduation curve of the integrated system of the piezoelectric accelerometer and the charge amplifier finally is:







At any rate, piezoelectric accelerometers have a determined *time constant* and therefore are *NOT* suited to measure a constant acceleration ($\omega = 0$) !

Piezoelectric accelerometers are factory calibrated one by one !

Manufacturer provide the user with the *transducer* sensitivity in [pC/ms⁻²] and the frequency response chart